# Cosmic Strings 

## some new results

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## Plan <br> 3 projects

- Quantum formation of topological defects.
- Evolution of global string loops.
- Evolution of gauge string loops.


## Examples of topological defects

Nematic liquid crystals and superconductors


Superconductor

## Emergence of classical structures from the quantum vacuum

## A full quantum calculation (kinks first)



Mainak Mukhopadhyay, George Zahariade \& TV, 2004.07249, 2009.11480

## A quantum mechanics problem

## A toy


but now in quantum field theory:

kinks ~ configurations trapped on top of V

## Formation

## the relevant physics




Solve the functional Schrodinger equation (for the ground state):

$$
P\left[\left\{\phi_{i}\right\}, t\right]=\Psi^{\dagger} \Psi=\frac{1}{\sqrt{\operatorname{det}(2 \pi K})} e^{-\phi^{T} K^{-1} \phi / 2}
$$

K is a time dependent $\mathrm{N} \times \mathrm{N}$ matrix.

$$
\begin{equation*}
K=Z Z^{\dagger} \quad \ddot{Z}+\Omega_{2}(t) Z=0 \tag{!}
\end{equation*}
$$

$Z$ is a time dependent, complex NxN matrix.

$$
\Omega_{2}(t)=-\nabla^{2}+m_{2}(t)
$$

## Formation

## counting kinks



Count zeros, i.e. sign changes.

$$
\begin{aligned}
& n_{\underline{Z}}=\left\langle\hat{n}_{Z}\right\rangle=\frac{N}{2 L}\left[1-\left\langle\operatorname{sgn}\left(\hat{\phi}_{1} \hat{\phi}_{2}\right)\right\rangle\right] \quad \text { Use translational invariance. } \\
& \left\langle\operatorname{sgn}\left(\hat{\phi}_{1} \hat{\phi}_{2}\right)\right\rangle=\frac{1}{\sqrt{\operatorname{det}(2 \pi K)}} \sum_{\underline{\text { quads. }}} \int d \phi_{1} \ldots d \phi_{N} \operatorname{sgn}\left(\phi_{1} \phi_{2}\right) e^{-\phi^{T} K^{-1} \phi / 2} \\
& n_{\underline{K}}=\frac{N}{2 L}\left[1-\frac{2}{\pi} \sin ^{-1}\left(\frac{\sum_{|n| \leq n_{c}}\left|c_{n}\right|^{2} \cos (2 \pi n / N)}{\sum_{|n| \leq n_{c}}\left|c_{n}\right|^{2}}\right)\right] \begin{array}{l}
\text { skipping quite a bit of } \\
\text { math.... }
\end{array}
\end{aligned}
$$

$$
\ddot{c}_{n}+\left[k_{n}^{2}+m_{2}(t)\right] c_{n}=0 \quad \text { with specified initial conditions. }
$$

## Formation

## results



Independent of $\boldsymbol{\tau}$ !
Compare: Kibble-Zurek

## Applicability

non-zero $\lambda$
Perturbation parameter: $\sim \lambda \tau / m$


Mainak Mukhopadhyay, George Zahariade \& TV, arXiv:2009.11480

## Formation

## vortices (2 spatial dimensions)

$$
\Phi=\phi+i \psi
$$

$$
L=\frac{1}{2}\left|\partial_{\mu} \Phi\right|^{2}-\frac{1}{2} m_{2}(t)|\Phi|^{2}-\frac{\lambda}{4}|\Phi|^{4}
$$

$$
m_{2}(t)=-m^{2} \tanh \left(\frac{t}{\tau}\right)
$$




## Formation

## vortices/strings



## Evolution

## Global strings (3 dimensions)

String core

$$
L=\frac{1}{2}\left|\partial_{\mu} \phi\right|^{2}+\frac{1}{2} m^{2}|\phi|^{2}-\frac{\lambda}{4}|\phi|^{4}
$$

Relevant to axion models before the QCD phase transition, where $\boldsymbol{\phi}$ is the Peccei-Quinn field. The phase of $\boldsymbol{\phi}$ is the axion field.

$$
\begin{aligned}
& \text { Straight string: } \phi=\eta f(r) e^{i \varphi} \\
& \qquad f(r) \sim " \tanh (r / w) "
\end{aligned}
$$

Goldstone cloud (extends to infinity)

Energy density falls of as $1 / r$ because of Goldstone cloud. Similar to electric line charge.

## Evolution

## Kalb-Ramond dynamics



Kalb-Ramond action in terms of 2-form field:

$$
\begin{gathered}
S=-\mu \int d^{2} \sigma \sqrt{-g_{2}}+\kappa \int d \sigma^{\mu \nu} A_{\mu \nu}-\frac{1}{6} \int d^{4} x H_{\mu \nu \lambda} H^{\mu \nu \lambda} \\
\text { Nambu-Goto } \quad \text { Goldstone cloud }
\end{gathered}
$$

Caveats: No massive radiation. Small backreaction.
Results: Goldstone boson radiation at primary frequency with $k \sim 1 / L$. Loop decays after $\sim 10$ oscillations. Tight constraints on QCD axion mass.

## Evolution

## Field theory dynamics

$$
\partial_{t}^{2} \phi_{a}=\nabla^{2} \phi_{a}-\frac{1}{2}\left(\phi_{b} \phi_{b}-1\right) \phi_{a} \quad a, b=1,2
$$

Caveats: Initial conditions? Limited by simulation size.

Results: Goldstone boson radiation with $1 / k$ power spectrum.
Loops decay within $\sim 1$ oscillation.
More relaxed constraints on QCD axion mass.
C. Hagmann \& P. Sikivie, 1991; T. Hiramatsu et al, 2011;
M. Gorghetto, E Hardy \& G. Villadoro, 2018; V.B. Klaer \& G. Moore, 2019

## Evolution

## Field theory simulations

## Parallel on XSEDE

What's a good way to set up the initial conditions?
Use straight string solution and mimic cosmological production of loops.


Technical note: Requires Lorentz boosting the static straight string solutions, patching together the four string solutions, and enforcing periodic boundary conditions. The latter requires modifications to the "product ansatz" for patching strings.
A. Saurabh, TV, \& L. Pogosian, 2020

## Animation

Total energy; potential energy

$$
|\mathbf{v}|=0.6
$$



## Results

## Core energy; angular momentum

 Core: $|\phi|<0.9 \eta$$$
L / w=(50,100,150,200,250) \times 4
$$



## Results

## Loop lifetime



## Results

Energy in massive and massless components
. $10^{4}$


## Results

Can we see the creation of massive modes?
Look for bound states in core.


$$
\begin{gathered}
\phi=\left(f(r)+e^{-i \omega t} g(r)\right) e^{i \theta} \\
-f^{\prime \prime}-\frac{f^{\prime}}{r}+\left[\frac{1}{r^{2}}-\frac{1}{2}\left(1-f^{2}\right)\right] f=0 \\
-g^{\prime \prime}-\frac{g^{\prime}}{r}+\left[\frac{1}{r^{2}}-\frac{3}{2}\left(1-f^{2}\right)\right] g=\Omega g
\end{gathered}
$$

$$
\Omega \equiv \omega^{2}-1=-0.19 \text { implies bound state. }
$$

Bound states excited by string intersections and Goldstone boson back reaction $\left(1 / r^{\wedge} 2\right.$ term $)$.

## Energy spectra

## Massive and massless modes



N=50 (blue), 100, 150, 200, 250.

Goldstone radiation


## Summary

## Global string loop results and caveats

We have simulated (cosmological) global string loops with length up to 1000 times the core width.

- Global string loops decay within ~1 oscillation period.
- Radiate massive and massless radiation according to initial energies.
- Massive particles are non-relativistic and eventually decay to massless radiation.
- Spectrum of massless radiation is $1 / k$.

Consistent with Hagmann \& Sikivie

- Caveat: Need to extrapolate by many orders of magnitude. Can't detect logarithmic effects.
(Code is available on request.)


## Evolution

## Gauge strings

## String core

$$
L=\frac{1}{2}\left|D_{\mu} \phi\right|^{2}-\frac{1}{4} F_{\mu \nu} F^{\mu \nu}+\frac{1}{2} m^{2}|\phi|^{2}-\frac{\lambda}{4}|\phi|^{4}
$$

Straight string: $\quad \phi=\eta f(r) e^{i \varphi}$

$$
A_{i}=v(r) \epsilon_{i j} \frac{x^{j}}{r^{2}}
$$

$$
f(r) \sim " \tanh (r / w) " \quad v(r) \sim " \tanh ^{2}(m r) "
$$

(No Goldstone cloud)
Energy density falls off as $\exp (-\mathrm{mr})$ because all fields are massive.

D. Matsunami, L. Pogosian, A. Saurabh, \& TV, 2020




## Evolution

## Gravitational waves or massive radiation?

Nambu-Goto action: $\quad S=-\mu \int d^{2} \sigma \sqrt{-g_{2}}$

Loops decay by gravitational radiation.
Full field theory simulations:
Loops decay by particle radiation.

TV \& A. Vilenkin, 1985; ...

Crucial to resolve for experiments (LIGO, NanoGrav,...) looking for gravitational wave signatures.

## Evolution

## Simulation equations

Technical note: Use Numerical Relativity technique for numerical stability.

$$
\begin{aligned}
\partial_{t}^{2} \phi_{a} & =\nabla^{2} \phi_{a}-e^{2} A_{i} A_{i} \phi_{a}-2 e \epsilon_{a b} \partial_{i} \phi_{b} A_{i}-e \epsilon_{a b} \phi_{b} \Gamma \\
& -\lambda\left(\phi_{b} \phi_{b}-\eta^{2}\right) \phi_{a} \\
\partial_{t} F_{0 i} & =\nabla^{2} A_{i}-\partial_{i} \Gamma+e\left(\epsilon_{a b} \phi_{a} \partial_{i} \phi_{b}+e A_{i} \phi_{a} \phi_{a}\right) \\
\partial_{t} \Gamma & =\partial_{i} F_{0 i}-g_{p}^{2} \frac{\left[\partial_{i} F_{0 i}+e \epsilon_{a b} \phi_{a} \partial_{t} \phi_{b}\right],}{\text { Gauss constraint }} \\
\Gamma & -\partial . \Lambda . \quad
\end{aligned}
$$

(Code is available on request.)

## Evolution

## Initial conditions



## Technical notes

Boost takes the gauge field out of temporal gauge. Then one needs to perform a gauge transformation to go back to temporal gauge.

Periodic boundary conditions require some smoothing functions.

## Evolution

## Animation



## Evolution

## Loop energy vs. time



## Evolution

## Lifetime vs. initial length


$\tau_{\text {particle }} \propto L^{2} \quad \tau_{\text {grav }} \propto L$

$$
\tau_{\text {grav }}<\tau_{\text {particle }} \text { for large } L
$$

$$
L_{\mathrm{crit}} \sim \frac{w}{G \mu}
$$

where $\mathrm{w}=$ width of the string, $\boldsymbol{\mu}=$ tension. Strings with tension above the QCD scale primarily decay by gravitational radiation.

High frequency cutoff on gravitational wave spectrum due to particle radiation.

## Gravitational radiation

## Kinks and cusps




FIG. 3: SBGW including the backreaction of particle emission on the loop distribution. LH panel: kinks on loops, RH panel: cusps on loop. The spectra are cutoff at high frequency, as indicated by the black vertical lines. $G \mu$ ranges from $10^{-17}$ (lower curve), through $10^{-15}, 10^{-13}, 10^{-11}, 10^{-9}$ and $10^{-7}$ (upper curve). Also plotted are the power-law integrated sensitivity curves from SKA (pink dashed) [44], LISA (yellow dashed) [45], adv-LIGO (grey dashed) [46] and Einstein Telescope (blue dashed) [47, 48].

## Conclusion

## Formation \& Evolution

- Formation: Universal results for the number density of topological defects formed in a quantum phase transition.
- Global string loop evolution: Loops decay in about 1 oscillation period, emit massive and massless Goldstone boson radiation. String core appears fluffy, probably due to excitation of bound states on core. Goldstone boson spectrum goes as $1 / k$ and with bump at $\mathrm{m}_{\chi} / 2$.
- Gauge string loop evolution: Loops larger than a critical length $w / G \mu$, decay primarily to gravitational radiation.

